

**CALCULATING THREE THERMAL COEFFICIENTS
FROM ONE DATA SET**

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Abstract: We study the problem of determining three thermal coefficients from one set data of a model problem rising in thermodynamics. This is an inverse problem, that is to coincide the solution of the differential equation with actual experimental results. The used method is based on minimizing the solution of the problem with the experimental data. Both the direct and inverse problems are described and numerical results are given.

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1. Introduction

The problem of finding the coefficients and parameters for thermodynamical problems when the solution is known is an inverse problem. To solve the inverse problem, one must first solve the direct problem, then solve the inverse problem for some coefficients and parameters. Solving such a problem therefore requires solving an optimization (minimization) problem, which is algorithmically more

challenging than the linear problem. These problems have many applications in engineering and science.

Inverse problems arise in many branches of science and engineering where the values of some sample material parameters must be obtained from the experimental data (see [10], [11], [12], [13], [14], [20], [21], [22]). Theory and applications of the determination of parameters of sample material has seen tremendous growth in recent years. Inverse problems can be formulated in many mathematical areas and can be analyzed by many different computational techniques (see [15], [16], [17], [18], [19], [23], [24], [25], [26]).

In [1] the authors illustrate the full description of fractal-based techniques and their application to the solution of inverse problems for ordinary and partial differential equations. In [2] the authors give a description of an inverse electromagnetic problem using a perturbation homotopy method combined with Gauss-Newton methods. In [3] the authors investigate a method for imposing two natural frequencies on a dynamic system consisting of an Euler-Bernoulli beam and carrying a single mass attachment. In [4] the authors use an algorithm to solve the problem of splicing the shredded paper. In [5] the authors describe an internal tidal model with experiments to investigate the estimation of spatially varying bottom friction coefficients.

The book by Evensen (2006) provides a good overview of many computational aspects of the subject, reflecting the author's experience in geophysical applications and related areas and provides a good entry point to some of the current research in this area. In [7] Kaipio and Somersalo introduce the Bayesian approach to inverse problems, especially in the context of differential equations. There is a wide research literature in the area of parameter estimation (see [8]), as well as attempts to introduce the notions of parameter estimation.

Wave propagation problems in environmental applications such as seismic analysis, acoustic and electromagnetic scattering are described in [9] for both forward and inverse problems.

In our thermodynamics model problem we will study the problem of determining three thermal coefficients from one set data. This is an inverse problem, that is to coincide the solution of the differential equation with actual experimental results.

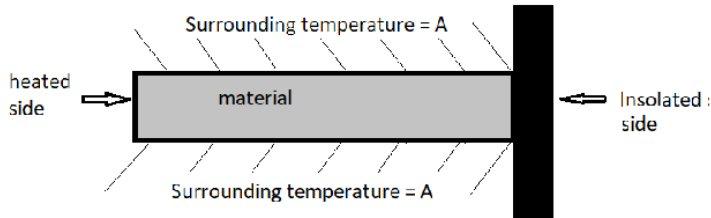


Figure 1: Geometry of the model problem.

2. The model problem

Our thermodynamic model problem consists on of a material that is insulated at $x = L$ and with heat source located at $x = 0$.

We assume the sample is initially at the uniform temperature, A , heated by a lamp at $x = 0$ (see Fig. 1).

The sample is assumed to be homogeneous solid material.

At the right end point is insulated, that is,

$$\frac{\partial u}{\partial x}(L, t) = 0.$$

At the left hand endpoint

$$W = -k \frac{\partial u}{\partial x}(0, t)$$

which means that the heat source is held constant during the course of the experiment.

Inside the sample, its temperature, $u(x, t)$ is governed by the Newton type heat equation. So, the problem of finding u involves solving

$$\begin{aligned} \rho c \frac{\partial u}{\partial t}(x, t) &= k \frac{\partial^2 u}{\partial x^2}(x, t) - m(u - A), \quad 0 < x < L, \quad t > 0, \\ u(x, 0) &= A, \quad 0 \leq x \leq L, \\ \frac{\partial u}{\partial x}(L, t) &= 0, \quad t > 0, \\ \frac{\partial u}{\partial x}(0, t) &= -W/k, \quad t > 0, \end{aligned} \quad (1)$$

where ρ is the density, c is the specific heat of the substance, m is the heat transfer coefficient, W is the power of the source, and k is the thermal conductivity. Here ρ is assumed known from an independent measurement. L is the length of the material.

The parameters k , c , and m are to be determined from measurements. It is simplest in dealing with the equation (1), to divide by ρc and set $a = k/(\rho c)$, the thermal diffusivity and $\alpha = m/(\rho c)$ so the partial differential equation is

$$\frac{\partial u}{\partial t}(x, t) = a \frac{\partial^2 u}{\partial x^2}(x, t) - \alpha(u - A). \quad (2)$$

We shall then determine parameters k , a , and α . If k , a , and α are known, then $c = k/(\rho a)$ and $m = \rho c \alpha$ are determined successively.

To simplify our problem, we make the change of variables to get the new problem for $v(x, t)$

$$\begin{aligned} \varphi(x) &= \frac{W}{2Lk} (x - L)^2 \\ u(x, t) &= v(x, t) + \varphi(x), \end{aligned} \quad (3)$$

$$\begin{aligned} v_t(x, t) &= a v_{xx}(x, t) - \alpha v(x, t) \\ +[a \varphi''(x) - \alpha \varphi(x) + \alpha A], \quad 0 < x < L, \quad t > 0 \end{aligned} \quad (4)$$

$$v_x(0, t) = 0, \quad v_x(L, t) = 0 \quad (5)$$

$$v(x, 0) = A - \varphi(x). \quad (6)$$

Once $v(x, t)$ is known, then $u(x, t)$ is obtained from (3).

To solve the problem for $v(x, t)$ we need to solve first the problem

$$z_t(x, t) = a z_{xx}(x, t) - \alpha z(x, t), \quad 0 < x < L, \quad t > 0 \quad (7)$$

$$z_x(0, t) = z_x(L, t) = 0 \quad (8)$$

$$z(x, 0) = A - \varphi(x) \quad (9)$$

which is done by setting $z(x, t) = X(x)T(t)$. The explicit solution is then formed by several steps to obtain the eigenfunctions.

$$z(x, t) = e^{-\alpha t} \left[\frac{WL}{6k} - \sum_{n=1}^{\infty} \frac{2W}{k\lambda_n^2 L} e^{-a\lambda_n^2 t} \cos(\lambda_n x) \right] \quad (10)$$

and

$$\begin{aligned} \zeta(x, t) &= \frac{4W}{\alpha k} \left(a + \frac{1}{3}L^2 \right) (1 - e^{-\alpha t}) \\ &- \sum_{n=1}^{\infty} \frac{2W}{kL\lambda_n^2 (\alpha + a\lambda_n^2)} \left[1 - e^{-(\alpha + a\lambda_n^2)t} \right] \cos(\lambda_n x), \end{aligned} \quad (11)$$

where $\lambda_n = n\pi/L$, $n = 1, 2, \dots$

With $z(x, t)$ and $\zeta(x, t)$ known, $u(x, t)$ is obtained from the equation

$$u(x, t) = z(x, t) + \zeta(x, t) + \varphi(x) + A. \quad (12)$$

3. The inverse problem

There are three parameters that must be identified: a , k , and α . They all appear in the equations. the physically relevant parameters are k , c , m and the relationship are

$$k \text{ is the same, } c = k/(\rho a), \quad m = \rho c \alpha,$$

where ρ is obtained independently.

Now the way the experiment is run as follows: Over a time period, $0 \leq t \leq T$, the lamp is turned on with a power Q watts. Then is turned off. The temperature then decreases until it reaches the ambient temperature along the full length of the sample. This actually means that the power of the lamp is a function of the time t so we should write

$$Q(t) = \begin{cases} W, & 0 \leq t \leq T, \\ 0, & t > T. \end{cases} \quad (13)$$

For $t > T$ one has to reset the problem where $u(x, T)$ is the “initial” condition and the boundary conditions are $u_x(0, t) = 0$ and $u_x(L, t) = 0$.

We need data at several points. The way the identification go is that we have a solution to the initial problem for all time. The time dependent part, that is the so called transient part tends to zero rapidly so choose T so large that the transient part of the solution is negligible. The time independent part of the solution satisfies:

$$\begin{aligned} ku_{xx}(x, t) - m[u(x, t) - A] &= 0 \\ u_x(0, t) &= -W/k \\ u_x(L, T) &= 0. \end{aligned}$$

The solution is

$$u(x) = \frac{W}{\sqrt{km}} \frac{\cosh(\sqrt{\frac{m}{k}}L)}{\sinh(\sqrt{\frac{m}{k}}L)} \cosh(\sqrt{\frac{m}{k}}x) - \sinh(\sqrt{\frac{m}{k}}x) + A. \quad (14)$$

The temperature gauges, that is, the thermal couples must be set in two separate places, say at $x = L/3$ and $x = 2L/3$.

One that solves for \sqrt{km} and $\sqrt{k/m}$ in terms of the measured values, say $u(L/3) = \mu$ and $u(2L/3) = \gamma$. To find a we go to the time dependent case.

Now k and m are known. a is chosen so that the measured value of the temperature, say $U(t)$, are as close as possible to the calculated values, that is,

$$\min_{\gamma} \left(\int_0^T [u(\frac{L}{3}, t) - U(t)]^2 dt \right)^{1/2} \quad (15)$$

is as small as possible. This will give the value for a .

4. Numerical results

In this section we will show some numerical results that determine the values of a , k and α and therefore $u(x, t)$ in the inverse problem. For the infinite series in equations (10) and (12) we took 100 terms to guarantee the convergence of the series. The integral in equation (15) is approximated using the trapezoidal rule. The least squares method along with Newton's method for nonlinear equation are used for the minimization of equation (15).

In this example we consider an unknown metal material with length $L = 5$ cm and density $\rho = 2.71$ g/cm³. The experiment was performed as described in Fig. 1 in the Laboratory of Zayed University, with the initial temperature of the sample $A = 20^0c.$, $W = 100$ watts. The experiment was run for 60 sec. and then the lamp was turned off to let the sample cool down. The results of the experiment is shown in Fig. 2 giving the measured values of $u(L/3, t)$ and $u(2L/3, t)$, $0 \leq t \leq 120$.

As $t \rightarrow \infty$, $u(L/3) \rightarrow 332.9^0c.$ and $u(2L/3) \rightarrow 322.6^0c$ (see Fig. 3) These values were used in Eqn. (14) to get the measured values for $k = 203$ W/cm-K and $m = 1.45$. These values are then used for the minimization of Eqn. (14) to obtain $a = 81.36$ cm/sec. and therefore $c = k/(\rho * a) = 0.92$ W-s/Kg-K.

Nomenclature:

k = conductivity (W/cm-K)

c = specific heat (W-s/Kg-K)

ρ = density (g/cm³)

W = power (W/cm²)

$a = k/\rho c$ thermal diffusivity (cm²/s)

A = initial temperature of sample (⁰ c.)

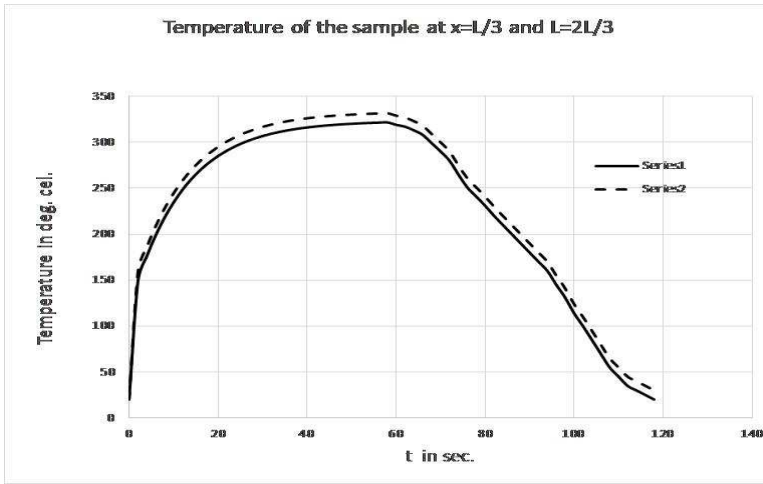


Figure 2: Experimental results of $u(x, t)$ at $x = L/3$ (series 2) and $x = 2L/3$ (series 1).

5. Conclusion

This paper deals with the determination of three thermal coefficients from one set data of a model problem arising in thermodynamics. The present study shows that we can easily get the thermal coefficients of a material by solving an inverse problem that leads to an optimization problem. The model problem was presented and numerical results were given.

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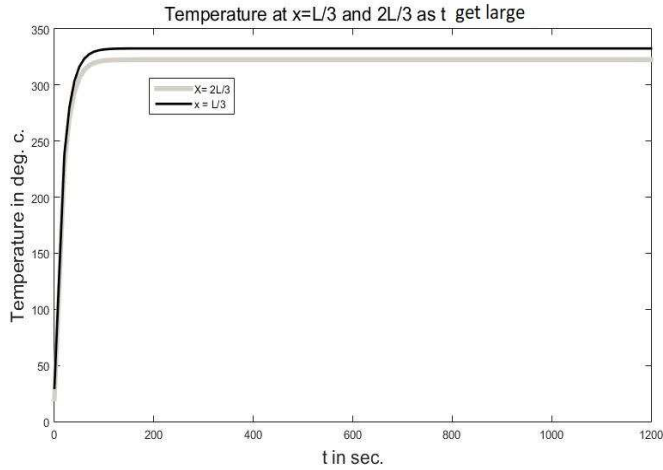


Figure 3: Temperature as t get large.

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