NUMERICAL SOLUTION OF THE TIME-FRACTIONAL
DIFFUSION EQUATIONS VIA QUARTER-SWEEP
PRECONDITIONED GAUSS-SEIDEL METHOD

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Abstract: In this research, we propose the approximate solution of the time-fractional diffusion equation based on a quarter-sweep implicit finite difference approximation equation. To derive this approximation equation, Caputo's time-fractional derivative has been used to discretize the proposed problems. By using the Caputo finite difference approximation equation, a linear system will be generated and solved iteratively. In addition to that, formulation and implementation the QSPGS iterative method are also presented. Based on the numerical results of the proposed iterative method, it can be concluded that the proposed iterative method is superior to the FSPGS and HSPGS iterative method.

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1. Introduction

In recent years, many studies use fractional partial differential equations (FPDEs) [14, 7, 15, 3, 2] for solving fractional problems to derive numerical and/or analytical solutions. Based on iterative methods for solving a one-dimensional diffusion model with constant coefficients and analytical solutions and for an instant a fractional derivative replaces the first-order space partial derivative in a diffusion model and lead to slower diffusion [14]. Therefore, there are numerical methods proposed to solve the time-fractional diffusion equations (TFDE), such as transform methods [22], finite elements together with the method of lines [15], explicit and implicit finite difference methods [15, 24]. Nevertheless, the explicit methods are conditionally stable, these finite difference schemes are available in the literature [24].

To solve the time-fractional diffusion equations (TFDE) problem needs to be discretized. Based on the implicit finite difference scheme and Caputo fractional operator, the approximation equations can be used to construct a linear system at each time level. To solve linear systems, many researchers also have discussed the concept of iterative methods, see [20, 10, 19] and reference therein. Besides these iterative methods, the concept of block iteration has also been introduced by [8]. Furthermore, Ibrahim and Abdullah [11], and Yousif and Evans [21] have pointed out the efficiency of block iterative methods.

For solving the large linear system, Abdullah [1] initiated Half-Sweep iteration, which is one of the most known and widely used iterative techniques to solve in solving any linear systems. Differently from the Half-sweep iteration approach, Othman and Abdullah [16] have expanded this approach to initiate the Modified Explicit Group (MEG) method based on the quarter-sweep approach. It is proved that this method is one of the most efficient block iterative methods in solving any linear system as compared with ED and EDG iterative methods. Also, another researcher has shown the capability of the quarter-sweep iteration in solving nonlinear system, see [5]. Among the existing iterative methods, the preconditioned iterative methods [4, 9] have been widely accepted to be one of the efficient methods for solving linear systems.

Because of the advantages of these iterative methods, this paper aims to construct and investigate the effectiveness of the Quarter-Sweep Preconditioned Gauss-Seidel (QSPGS) iterative method for solving time fractional parabolic partial differential equations (TPPDE's) based on the Caputo implicit finite difference approximation equation. In this paper, we investigate the performance of the Quarter-Sweep Preconditioned Gauss-Seidel (QSGS) iterative method for solving time-fractional parabolic partial differential equations (TPPDE's)

based on the Caputo implicit finite difference approximation equation. To demonstrate the capability of the Quarter-Sweep Preconditioned Gauss-Seidel (QSPGS) method, we also implement the Full-Sweep Preconditioned Gauss-Seidel or FSPGS and Half-Sweep Preconditioned Gauss-Seidel or HSPGS iterative methods being used as a control method.

2. Preliminaries

To begin the derivation of the QSPGS iterative method, let us consider the time-fractional diffusion equation (TFDE's) defined as

$$\frac{\partial^\alpha u(x, t)}{\partial^\alpha t} = a(x) \frac{\partial^2 u(x, t)}{\partial x^2} + b(x) \frac{\partial u(x, t)}{\partial x} + c(x) u(x, t), \quad (1)$$

where $a(x)$, $b(x)$ and $c(x)$ are known functions or constants whereas α is a parameter which refers to the fractional-order of time derivative. Before to set the discretizing problem (1), let us remind some definitions from the theory of fractional calculus.

Definition 1. ([17]) The Riemann-Liouville integral operator J^α of fractional order α is defined as

$$J^\alpha f(x) = \frac{1}{\Gamma(\alpha)} \int_0^x (x-t)^{\alpha-1} f(t) dt, \quad \alpha > 0, x > 0. \quad (2)$$

Definition 2. ([17]) The Caputo fractional derivative operator D^α of order $\alpha > 0$ is defined as

$$D^\alpha f(x) = \frac{1}{\Gamma(m-\alpha)} \int_0^x \frac{f^{(m)}(t)}{(x-t)^{\alpha-m+1}} dt, \quad (3)$$

with $m-1 < \alpha \leq m$, $m \in \mathbb{N}$, $x > 0$. In (2) and (3), $\Gamma(\alpha)$ is the well-known Gamma function

$$\Gamma(\alpha) = \int_0^\infty x^{\alpha-1} e^{-x} dx. \quad (4)$$

For solving the numerical of time-fractional diffusion equation (TFDE's), in equation (1), we get numerical approximations by using the Caputo derivative definition with Dirichlet boundary conditions and consider the non-local fractional derivative operator. This approximation equation can be categorized as an unconditionally stable scheme. On strength of Problem (1), the solution

domain of the problem has been restricted to the finite space domain $0 \leq x \leq \gamma$, with $0 < \alpha < 1$ whereas the parameter α refers to the fractional-order of space derivative. To solve Problem (1), let us consider the initial and boundary conditions of Problem (1) be given as

$$u(0, t) = g_0(t), u(l, t) = g_1(t), \quad (5)$$

and the initial condition

$$u(x, 0) = f(x), \quad (6)$$

where $g_0(t)$, $g_1(t)$, and $f(x)$ are given functions. Based on a discretized approximation to the time-fractional derivative in equation (1), we consider Caputo's fractional partial derivative of order α , as

$$\frac{\partial^\alpha u(x, t)}{\partial^\alpha t} = \frac{1}{\Gamma(n-1)} \int_0^\infty \frac{\partial u(x-s)}{\partial t} (t-s)^{-\alpha} ds, \quad t > 0, 0 < \alpha < 1. \quad (7)$$

3. The Caputo implicit finite difference approximation

Based on equation (7), the formulation of the Caputo fractional partial derivative of the first order approximation method is given as

$$D_t^\alpha U_{i,n} \cong \sigma_{\alpha,k} \sum_{j=1}^n \omega_j^\alpha (U_{i,n-j+1} - U_{i,n-j}), \quad (8)$$

and we have the following expressions

$$\sigma_{\alpha,k} = \frac{1}{\Gamma(1-\alpha)(1-\alpha)^{1-\alpha}}, \quad (9)$$

and

$$\omega_j^\alpha = j^{1-\alpha} - (j-1)^{1-\alpha}. \quad (10)$$

First, to discretize Problem (1), let the solution domain of the problem be partitioned uniformly. To do this, we consider some positive integers m and n in which the grid sizes in space and time directions for the finite-difference algorithm are defined as $h = \delta x = \gamma/m$ and $k = \delta t = T/n$, respectively. According to these grid sizes, we develop the uniform grid network of the solution domain where the grid points in the space interval $[0, \gamma]$ are shown as the numbers $x_i = ih, i = 0, 1, 2, \dots, m$, and the grid points in the time interval $[0, T]$ are labeled $t_j = jk, j = 0, 1, 2, \dots, n$. Then the values of the function $U(x, t)$ at the grid points are denoted as $U_{i,j} = U(x_i, t_j)$. According to equation (8) and the implicit finite difference discretization scheme, the Caputo implicit finite

difference approximation equation of Problem (1) to the grid point centered at $(x_i, t_j) = (ih, jk)$ is given as

$$\sigma_{\alpha,k} \sum_{j=1}^n \omega_j^\alpha (U_{i,n-j+1} - U_{i,n-j}) = a_i \frac{1}{16h^2} (U_{i-4,n} - 2U_{i,n} + U_{i+4,n}) + b_i \frac{1}{8h} (U_{i+4,n} - U_{i-4,n}) + c_i U_{i,n}, \quad (11)$$

for $i = 4, 8, \dots, m-4$. Thus, based on equation (11), this approximation equation is known as the fully implicit finite difference approximation equation which is consistent first-order accuracy in time and second-order in space. Particularly, the approximation equation (11) can be rewritten based on the specified time level. Immediately, we have for $n \geq 2$:

$$\sigma_{\alpha,k} \sum_{j=1}^n \omega_j^\alpha (U_{i,n-j+1} - U_{i,n-j}) = p_i U_{i-4,n} + q_i U_{i,n} + r_i U_{i+4,n}, \quad (12)$$

where

$$p_i = \frac{a_i}{16h^2} - \frac{b_i}{8h}, q_i = c_i - \frac{a_i}{8h^2}, r_i = \frac{a_i}{16h^2} + \frac{b_i}{8h}. \quad (13)$$

Also, we get for $n = 1$,

$$-p_i U_{i-4,1} + q_i U_{i,1} - r_i U_{i+4,1} = f_{i,1}, i = 4, 8, \dots, m-4, \quad (14)$$

where $\omega_j^{(\alpha)} = 1, q_i^* = \sigma_{\alpha,k} - q_i, f_{i,1} = \sigma_{\alpha,k} U_{i,1}$. Furthermore, based on equation (14), it can be seen that the tridiagonal linear system can be constructed in matrix form as

$$AU = f, \quad (15)$$

where

$$A = \begin{bmatrix} q^*_4 & -r_4 & & & & \\ -p_8 & q^*_8 & -r_8 & & & \\ & -p_{12} & q^*_{12} & -r_{12} & & \\ & & \ddots & \ddots & \ddots & \\ & & & -p_{m-8} & q^*_{m-8} & -r_{m-8} \\ & & & & -p_{m-4} & q^*_{m-4} \end{bmatrix}_{(\frac{m}{4}-1) \times (\frac{m}{4}-1)}, \quad (16)$$

$$U = [U_{4,1} \ U_{8,1} \ U_{12,1} \ \dots \ U_{m-8,1} \ U_{m-4,1}], \quad (17)$$

$$f = [U_{4,1} + p_1 U_{0,1} \ U_{8,1} \ U_{12,1} \ \dots \ U_{m-8,1} \ U_{m-4,1} + p_{m-4} U_{m,1}]. \quad (18)$$

4. Analysis of stability

In this section, we have considered the stability analysis of the implicit finite difference approximation equation in equation (11). For stability analysis, we will use Von-Neumann's [13] and the Lax equivalence theorem [18]. It follows that the numerical solution of the approximation equation in equation (11) converges to the exact solution as $h, k \rightarrow 0$.

Theorem 3. *The fully implicit numerical method (11), the solution to equation (1) with $0 < \alpha < 1$ on the finite domain $0 \leq x \leq 1$, with zero boundary condition $u(0, t) = u(1, t) = 0$ for all $t \geq 0$ is consistent and unconditionally stable.*

Proof. To examine the stability of the proposed method, we find for the solution of the form $U_j^n = \zeta_n e^{i\omega jh}$, $i = \sqrt{-1}$, ω is real. Therefore, equation (12) becomes

$$\begin{aligned} \sigma_{\alpha,k} \zeta_{n-1} e^{i\omega jh} - \sigma_{\alpha,k} \sum_{j=2}^n \omega_j^{(\alpha)} (\zeta_{n-j+1} e^{i\omega jh} - \zeta_{n-j} e^{i\omega jh}) \\ = -p_i \zeta_n e^{i\omega(j-4)h} + (\sigma_{\alpha,k} - q_i) \zeta_n e^{i\omega jh} - r_i \zeta_n e^{i\omega(j+4)h}. \end{aligned} \quad (19)$$

By simplifying and reordering over equation (19), we have

$$\begin{aligned} \sigma_{\alpha,k} \zeta_{n-1} - \sigma_{\alpha,k} \sum_{j=2}^n \omega_j^{(\alpha)} (\zeta_{n-j+1} - \zeta_{n-j}) \\ = \zeta_n ((-p_i - r_i) \cos(\omega h)) + (\sigma_{\alpha,k} - q_i). \end{aligned} \quad (20)$$

This can be reduced to

$$\zeta_n = \frac{\zeta_{n-1} + \sum_{j=2}^n \omega_j^{(\alpha)} (\zeta_{n-j} - \zeta_{n-j+1})}{1 + \frac{(-p_i - r_i)}{\sigma_{\alpha,k}} \cos(\omega h) - \frac{q_i}{\sigma_{\alpha,k}}}. \quad (21)$$

From equation (21), it can be observed that the conducted as

$$\left(1 + \frac{(-p_i - r_i)}{\sigma_{\alpha,k}} \cos(\omega h) - \frac{q_i}{\sigma_{\alpha,k}}\right) \geq 1, \quad (22)$$

for all α, n, ω, h and k , we have

$$\zeta_1 \leq \zeta_0, \quad (23)$$

and

$$\zeta_n \leq \zeta_{n-1} + \sum_{j=2}^n \omega_j^{(\alpha)} (\zeta_{n-j} - \zeta_{n-j+1}), \quad n \geq 2. \quad (24)$$

Thus, for $n = 2$, the last inequality implies

$$\zeta_2 \leq \zeta_1 + \omega_2^{(\alpha)}(\zeta_0 - \zeta_1). \quad (25)$$

Again, repeating the above process, we can get

$$\zeta_j \leq \zeta_{j-1}, \quad j = 1, 2, \dots, n-1. \quad (26)$$

From equation (24), we finally have

$$\zeta_n \leq \zeta_{n-1} + \sum_{j=2}^n \omega_j^{(\alpha)}(\zeta_{n-j} - \zeta_{n-j+1}) \leq \zeta_{n-j}. \quad (27)$$

Since each term in the summation is negative, it shows that the inequalities (23) and (24) imply

$$\zeta_n \leq \zeta_{n-1} \leq \zeta_{n-2} \leq \dots \leq \zeta_1 \leq \zeta_0. \quad (28)$$

Thus

$$\zeta_n = |U_j^n| \leq \zeta_0 = |U_j^0| = |f_j|, \quad (29)$$

which entails $\|U_j^n\| \leq \|f_j\|$, and we have stability. $\square$

5. Formulation of Quarter-Sweep Preconditioned Gauss-Seidel (QSPGS)

Concerning the tridiagonal linear system (16), the characteristics of its coefficient matrix are large scale and sparse. As mentioned in Section 1, many researchers have discussed various iterative methods, such as Gunawardena [9] and Young [20]. To obtain numerical solutions of the tridiagonal linear system (16), we consider the Quarter-Sweep Preconditioned Gauss-Seidel (QSPGS) iterative method, which is the most known and widely used for solving any linear system. Before applying the QSPGS iterative method, we need to transform the original linear system (15) into the preconditioned linear system

$$A^*x = f^*, \quad (30)$$

where $A^* = PAP^T$, $f^* = Pf$, $U = P^Tx$. The matrix P is called a preconditioned matrix and defined as [12]

$$P = I + S, \quad (31)$$

where

$$S = \begin{bmatrix} 0 & -r_1 & 0 & 0 & 0 & 0 \\ 0 & 0 & -r_2 & 0 & 0 & 0 \\ 0 & 0 & 0 & -r_3 & 0 & 0 \\ 0 & 0 & 0 & 0 & \ddots & 0 \\ 0 & 0 & 0 & 0 & 0 & -r_{m-1} \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}_{(m-1) \times (m-1)}, \quad (32)$$

and the matrix I is an identical matrix. To formulate the QSPGS method, let the coefficient matrix A^* in (30) be expressed as the summation of the three matrices

$$A^* = D - L - V, \quad (33)$$

where D , L , and V are diagonal, lower triangular, and upper triangular matrices, respectively. By using equation (30), the formulation of QSPGS iterative method can be defined generally as

$$x^{(k+1)} = (D - L)^{-1} V x^{(k)} (D - L)^{-1} f^*, \quad (34)$$

where $x^{(k+1)}$ represents an unknown vector at $(k + 1)$ th iterations. The implementation of the QSPGS iterative method can be described in Algorithm 1.

Algorithm 1: QSPGS

i. Initialize $U \leftarrow 0$ and $\epsilon \leftarrow 10^{-10}$,

ii. For $j = 1, 2, \dots, n$, implement

For $i = 1, 2, \dots, m - 1$, calculate
 $x^{(k+1)} = (D - L)^{-1} V x^{(k)} + (D - L)^{-1} f^*$

Convergence test. If the convergence criterion i.e. $\|U^{(k+1)} - U^{(k)}\| \leq \epsilon$ is satisfied, go to Step (iii). Otherwise, go back to Step (i).

iii. Display approximate solutions.

6. Numerical experiment

In this section, we use one example of the time-fractional diffusion equation to show the accuracy and effectiveness properties of the Quarter Sweep Preconditioned Gauss-Seidel (QSPGS) compare Full-Sweep Preconditioned Gauss-Seidel (FSPGS) and Half Sweep Preconditioned Gauss-Seidel (HSPGS) iterative methods. These three parameters were executed on the computer using a program written in *C language*. For comparison purpose, three parameters will

be considered such as the number of iterations, execution time (in seconds), and maximum absolute error at three different values of $\alpha = 0.25, 0.50$, and 0.75 . For the implementation of these three iterative schemes, the convergence test considered the tolerance error, which is fixed as $\epsilon = 10^{-10}$. To illustrate the performance of QSPGS iteration method, let us consider the time-fractional initial boundary value problem be given as follows.

Example 1: ([6])

$$\frac{\partial^\alpha u(x, t)}{\partial t^\alpha} = \frac{\partial^2 u(x, t)}{\partial x^2}, \quad 0 < \alpha \leq 1, \quad 0 \leq x \leq \gamma, \quad t > 0, \quad (35)$$

where the boundary conditions are stated in fractional terms

$$u(0, t) = \frac{2kt^\alpha}{\Gamma(\alpha + 1)}, \quad u(l, t) = l^2 + \frac{2kt^\alpha}{\Gamma(\alpha + 1)}, \quad (36)$$

and the initial condition $u(x, 0) = x^2$. Following Problem (35), as taking $\alpha = 1$, it can be seen that equation (35) can be reduced to the standard diffusion equation

$$\frac{\partial u(x, t)}{\partial t} = \frac{\partial^2 u(x, t)}{\partial x^2}, \quad 0 \leq x \leq \gamma, \quad t > 0, \quad (37)$$

subject to the initial condition $u(x, 0) = x^2$, and the boundary conditions $u(0, t) = 2kt, u(l, t) = l^2 + 2kt$. Then, the analytical solution of equation (37) is obtained as follows:

$$u(x, t) = x^2 + 2kt. \quad (38)$$

Now, by applying the series

$$u(x, t) = \sum_{n=0}^{m-1} \frac{\partial^n u(x, 0)}{\partial t^n} \frac{t^n}{n!} + \sum_{n=1}^{\infty} \sum_{n=0}^{m-1} \frac{\partial^{mn+i} u(x, 0)}{\partial t^{mn+i}} \frac{t^{n\alpha+i}}{\Gamma(n\alpha+i+1)}, \quad (39)$$

$u(x, t)$ for $0 < \alpha \leq 1$, it can be shown that the analytical solution of equation (35) is given as

$$u(x, t) = x^2 + 2k \frac{t^\alpha}{\Gamma(\alpha + 1)}. \quad (40)$$

Example 2: ([6]) Let us consider the following time-fractional initial boundary value problem be defined as

$$\frac{\partial^\alpha u(x, t)}{\partial t^\alpha} = \frac{1}{2} x^2 \frac{\partial^2 u(x, t)}{\partial x^2}, \quad 0 < \alpha \leq 1, \quad 0 \leq x \leq \gamma, \quad t > 0, \quad (41)$$

where the boundary conditions are given as $u(0, t) = 0, u(1, t) = e^t$ and the initial condition $u(x, 0) = x^2$. From equation (41), as taking $\alpha = 1$, it can be shown that equation (41) can also be reduced to the standard diffusion equation

$$\frac{\partial u(x, t)}{\partial t} = \frac{1}{2} x^2 \frac{\partial^2 u(x, t)}{\partial x^2}, \quad 0 \leq x \leq \gamma, \quad t > 0. \quad (42)$$

Then, the analytical solution of equation (42) is obtained as follows:

$$u(x, t) = x^2 e^t. \quad (43)$$

Now, by applying the series

$$u(x, t) = \sum_{n=0}^{m-1} \frac{\partial^n u(x, 0)}{\partial t^n} \frac{t^n}{n!} + \sum_{n=1}^{\infty} \sum_{n=0}^{m-1} \frac{\partial^{mn+i} u(x, 0)}{\partial t^{mn+i}} \frac{t^{n\alpha+i}}{\Gamma(n\alpha+i+1)}, \quad (44)$$

$u(x, t)$ for $0 < \alpha \leq 1$, it can be shown that the analytical solution of equation (41) is stated as

$$u(x, t) = x^2 \left[1 + \frac{t^\alpha}{\Gamma(\alpha+1)} + \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} + \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} + \dots \right]. \quad (45)$$

All results of numerical experiments for equation (35) and (41), which were obtained from the implementation of FSPGS, HSPGS and QSPGS iterative methods have been recorded in Table 1 until Table 6 at different values of mesh sizes, $m = 128, 256, 512, 1024$, and 2048 .

7. Conclusion

As a conclusion for the numerical solution of the time-fractional diffusion problems, this paper deals with the implementation of the QSPGS iterative method to solve a linear system generated by the Quarter-Sweep Caputo implicit approximation equations. Through numerical experiments results from such Tables 1, 2 and 3 by comparing the performance between the FSPGS, HSPGS and QSPGS iterative methods at three different values of $\alpha = 0.25, 0.50$ and 0.75 , it can be seen that the percentage reduction of the number of iterations for the QSGS iterative method has declined approximately by $71.99 - 92.72\%$, $50.71 - 95.53\%$, and $40.41 - 96.45\%$ respectively as compared with the FSPGS and HSPGS method. Implementations of computational time for the QSPGS method are much faster about $69.78 - 98.09\%$, $30.64 - 88.35\%$, and $25.03 - 89.99\%$, respectively, than the FSPGS and HSPGS method. It can be concluded that the QSPGS method involves less number of iterations and computational time as compared with FSPGS and HSPGS methods. According to the accuracy of FSPGS, HSPGS, and QSPGS iterative methods, it can be stated that the numerical solutions of both methods are in good agreement.

Table 1: Numerical result of Example 1 at $\alpha = 0.25$

M	Method	k	time	Max Error
128	FSPGS	7292	35.86	9.96e-05
	HSPGS	1966	5.64	9.96e-05
	QSPGS	528	1.41	9.96e-05
256	FSPGS	26884	261.56	9.98e-05
	HSPGS	7292	37.36	9.96e-05
	QSPGS	1966	5.45	9.96e-05
512	FSPGS	98422	1916.28	1.00e-04
	HSPGS	26884	272.45	9.98e-05
	QSPGS	7292	35.88	9.96e-05
1024	FSPGS	357258	14064.44	1.04e-04
	HSPGS	98422	2025.13	1.04e-04
	QSPGS	26884	267.44	9.98e-05
2048	FSPGS	1183293	4104.17	1.36e-04
	HSPGS	339197	3121.13	1.36e-04
	QSPGS	94141	1120.08	9.98e-05

Table 2: Numerical result of Example 1 at $\alpha = 0.50$

M	Method	k	time	Max Error
128	FSPGS	4715	2.23	9.84e-05
	HSPGS	1270	1.59	9.84e-05
	QSPGS	342	0.98	9.84e-05
256	FSPGS	17417	16.68	9.87e-05
	HSPGS	4715	10.28	9.87e-05
	QSPGS	1270	3.78	9.87e-05
512	FSPGS	63298	123.01	9.96e-05
	HSPGS	17417	95.09	9.87e-05
	QSPGS	4715	23.31	9.87e-05
1024	FSPGS	232784	1007.47	1.03e-05
	HSPGS	63928	893.24	9.95e-05
	QSPGS	17417	171.84	9.87e-05
2048	FSPGS	1150153	3239.84	1.34e-05
	HSPGS	232784	2511.66	1.34e-05
	QSPGS	61246	975.43	9.98e-05

Table 3: Numerical result of Example 1 at $\alpha = 0.75$

M	Method	k	time	Max Error
128	FSPGS	2319	1.93	1.30e-04
	HSPGS	625	1.03	1.30e-04
	QSPGS	170	0.87	1.30e-04
256	FSPGS	8585	12.37	1.30e-04
	HSPGS	2319	8.23	1.30e-04
	QSPGS	625	2.27	1.30e-04
512	FSPGS	31619	62.78	1.31e-04
	HSPGS	8585	46.08	1.31e-04
	QSPGS	2319	11.80	1.31e-04
1024	FSPGS	115617	820.93	1.35e-04
	HSPGS	31619	636.78	1.35e-04
	QSPGS	8585	84.30	1.35e-04
2048	FSPGS	362784	1305.50	1.35e-04
	HSPGS	115627	807.13	1.35e-04
	QSPGS	31691	629.00	1.35e-04

Table 4: Numerical result of Example 2 at $\alpha = 0.25$

M	Method	k	time	Max Error
128	FSPGS	2873	8.48	1.95e-02
	HSPGS	774	5.05	1.95e-02
	QSPGS	209	3.42	1.95e-02
256	FSPGS	10624	96.54	1.95e-02
	HSPGS	2873	19.09	1.95e-02
	QSPGS	774	6.93	1.95e-02
512	FSPGS	39608	648.25	1.95e-02
	HSPGS	10624	108.69	1.95e-02
	QSPGS	2873	18.58	1.95e-02
1024	FSPGS	142635	791.55	1.95e-02
	HSPGS	39068	582.43	1.95e-02
	QSPGS	10624	107.34	1.95e-02
2048	FSPGS	487355	2543.23	1.95e-02
	HSPGS	128676	1326.21	1.95e-02
	QSPGS	36470	275.38	1.95e-02

Table 5: Numerical result of Example 2 at $\alpha = 0.50$

M	Method	k	time	Max Error
128	FSPGS	1398	7.00	8.28e-02
	HSPGS	378	5.11	8.28e-02
	QSPGS	104	3.29	8.28e-02
256	FSPGS	5162	35.69	8.29e-02
	HSPGS	1398	12.12	8.29e-02
	QSPGS	378	6.06	8.28e-02
512	FSPGS	18957	277.23	8.29e-02
	HSPGS	5162	57.19	8.29e-02
	QSPGS	1398	11.90	8.29e-02
1024	FSPGS	69108	492.97	8.29e-02
	HSPGS	38957	390.80	8.30e-02
	QSPGS	5162	56.40	8.30e-02
2048	FSPGS	240051	1781.32	8.29e-02
	HSPGS	67817	920.14	8.30e-02
	QSPGS	19430	252.12	8.28e-02

Table 6: Numerical result of Example 2 at $\alpha = 0.75$

M	Method	k	time	Max Error
128	FSPGS	655	4.44	1.37e-01
	HSPGS	178	2.59	1.37e-01
	QSPGS	50	1.20	1.37e-01
256	FSPGS	2420	15.95	1.37e-01
	HSPGS	655	8.50	1.37e-01
	QSPGS	178	5.56	1.37e-01
512	FSPGS	8911	184.75	1.37e-01
	HSPGS	2420	30.08	1.37e-01
	QSPGS	655	8.36	1.37e-01
1024	FSPGS	32602	420.11	1.37e-01
	HSPGS	8911	189.50	1.37e-01
	QSPGS	2420	29.49	1.37e-01
2048	FSPGS	116801	951.53	1.37e-01
	HSPGS	33318	511.32	1.37e-01
	QSPGS	8911	188.66	1.37e-01

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